

The vacuum comes alive

James Binney

Rudolf Peierls Centre for
Theoretical Physics

Outline

- From action at a distance to a dynamic aether
- The fundamental symmetry of the aether:
Lorentz invariance
- The language of field theory: Lagrangians
- How Lorentz invariance of the Lagrangian
largely specifies the dynamics

Newton

1642 - 1727

- Gravity: action at a distance

$$F \propto \frac{Mm}{r^2}$$



Coulomb

1736 - 1806

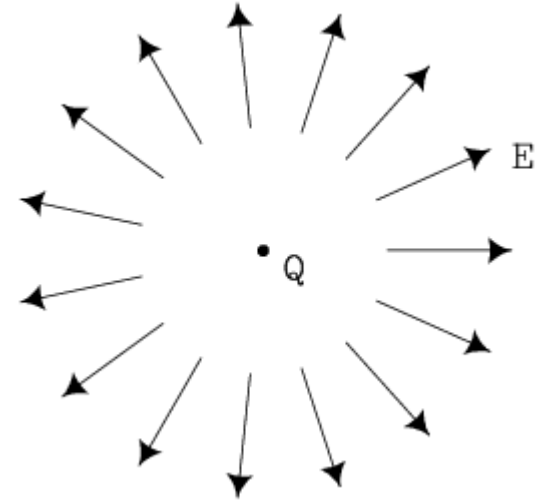


- Electrostatics

$$F \propto \frac{Qq}{r^2}$$

- Electric field as an abstraction

$$F = qE \quad E(r) \propto \frac{Q}{r^2}$$



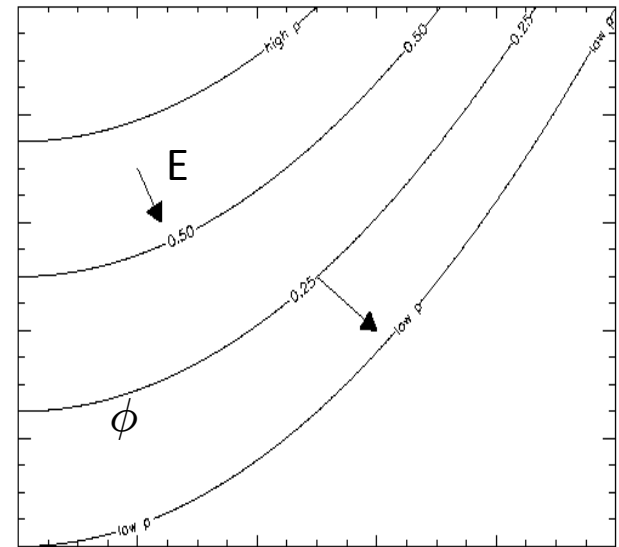
- Charges make a field; field pulls charges
- Same idea applies to gravity: $F = mg$

charge field

Potentials

- These fields are gradients (steepness) of a potential

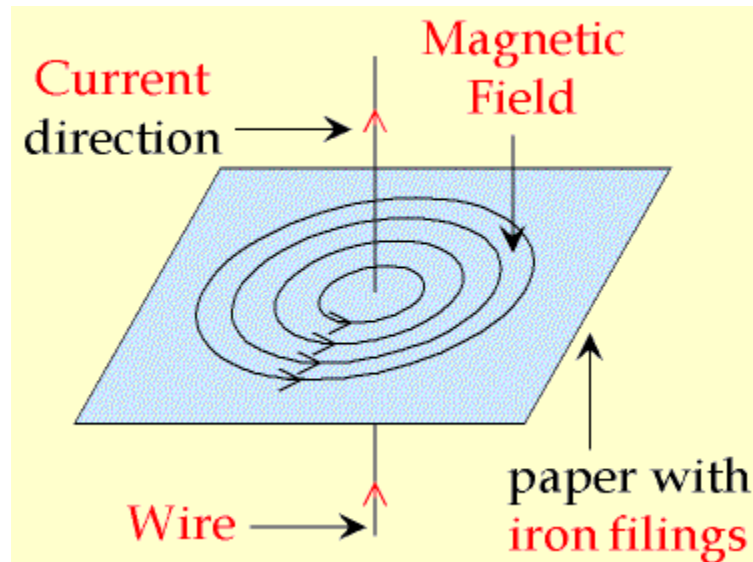
$$E = -\nabla\phi \quad \phi = \frac{\kappa_E Q}{r}$$
$$g = -\nabla\Phi \quad \Phi = -\frac{\kappa_G M}{r}$$



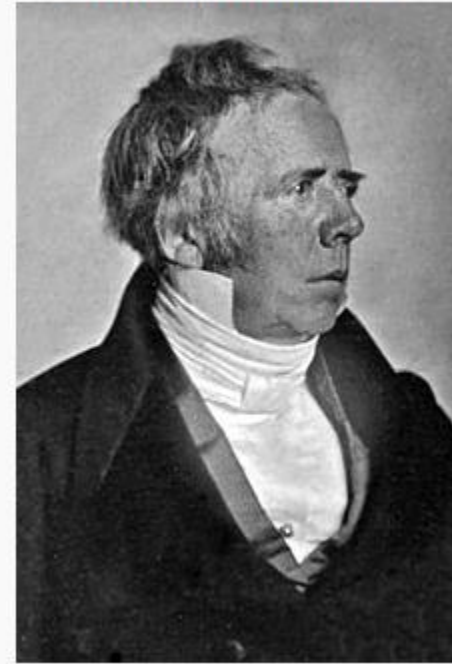
- Potential easier to work with because just one number ϕ versus E_x , E_y , E_z

Oersted (1777-1851)

- Current flowing in a wire causes a B field to run around the wire so compass needle always points perpendicular to compass—wire line



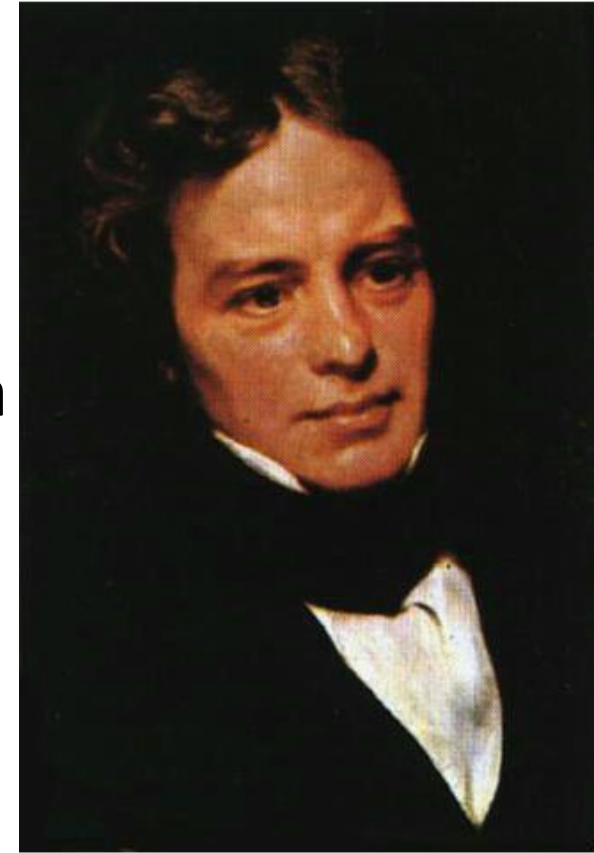
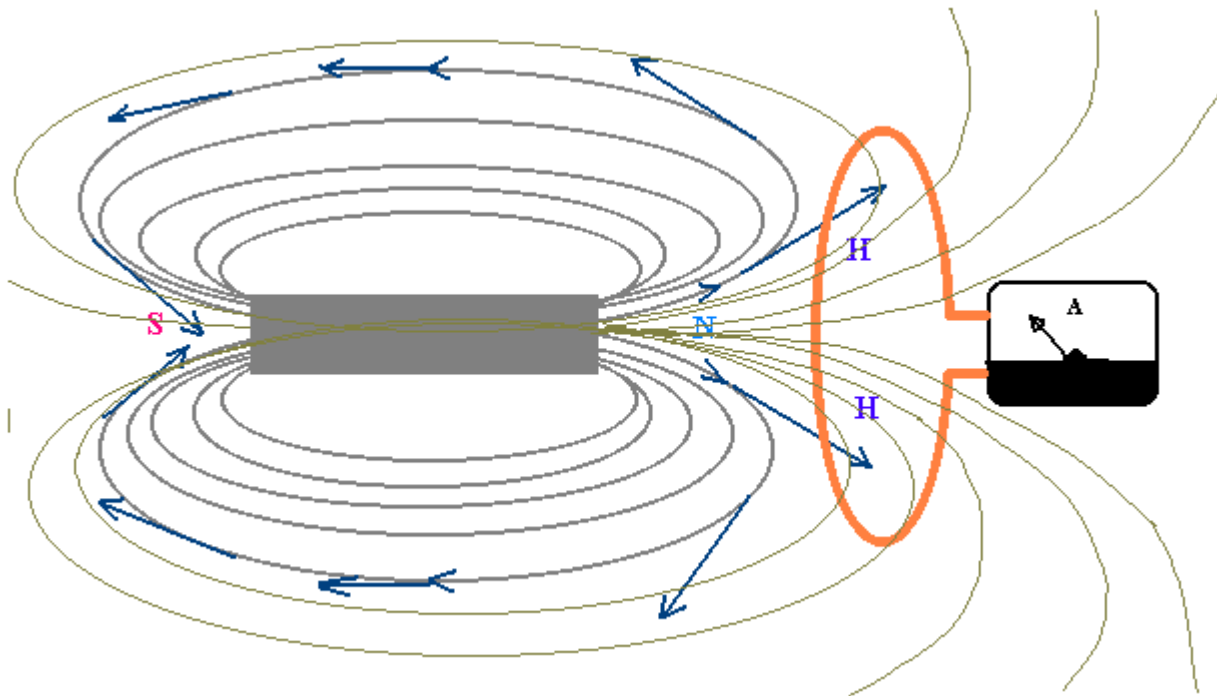
Hans Christian Ørsted



Faraday

1791 - 1867

- A changing magnetic field generates a swirling electric field

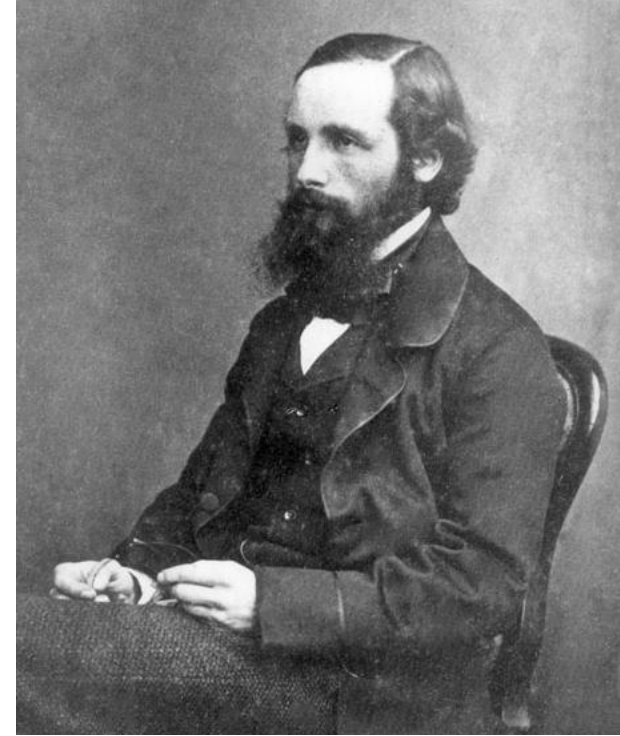


$$\nabla \times E = -\frac{\partial B}{\partial t}$$

Maxwell

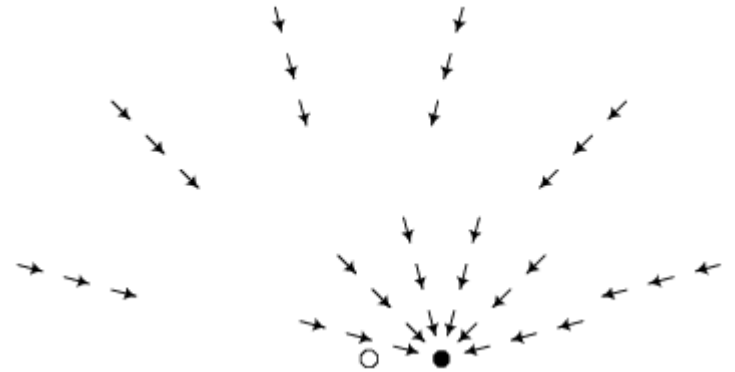
1831 - 1879

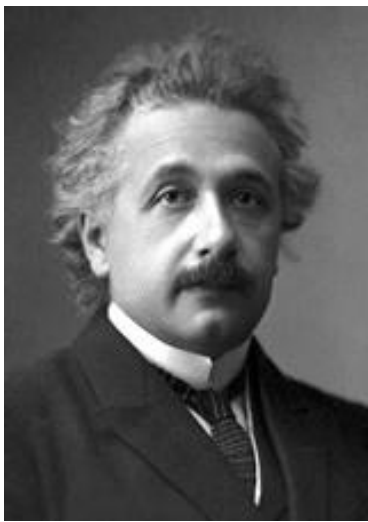
- A changing electric field generates a swirling magnetic field
- \Rightarrow E, B form an independent dynamical system
- Don't need charges to generate them!
- Ripples in E,B travel at speed of light



Maxwell (2)

- EB are manifestations of an “elastic” medium – the aether
- Action isn't @ a distance!
- And it isn't instantaneous
- A kink in E travels through the aether @ c
- An oscillating electric charge must radiate electromagnetic waves





Lorentz & Einstein

1853 – 1928 1879 - 1955



- Lorentz:

- Maxwell's equations unchanged under

$$\begin{pmatrix} ct' \\ x' \end{pmatrix} = \begin{pmatrix} \gamma & \beta\gamma \\ \beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} ct \\ x \end{pmatrix} \quad \text{where } \gamma^2 = \frac{1}{1 - \beta^2}$$

- What does this mean?

- Einstein:

- equation relates times and positions measured by moving observers ($\beta = v/c$)

- Invariance of equations implies physics is the same no matter how fast you move

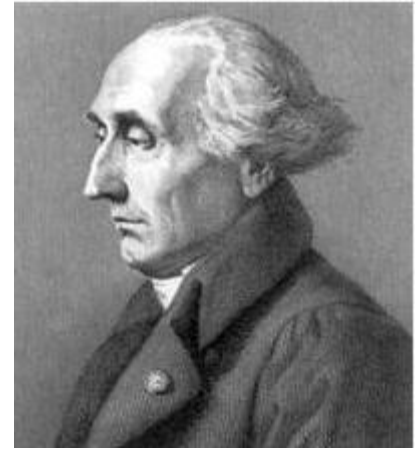
20thc physics

- Aether = spacetime
- The material world is nothing but the excited aether
- Fundamental equations of physics must be invariant under Lorentz's transformation
- This requirement restricts the dynamics of the aether to just a handful of possibilities
- Lagrangian densities enable us to find these possibilities

Lagrange

Giuseppe Luigi Lagrangia

1736 1813



- Projectiles etc follow path of “least action” $S = \int dt L(x,v)$
- $L = KE - PE$
- Given $L(x,v)$ we can derive the eqns of motion from the “Euler-Lagrange” eqns
- A field evolves so as to minimise action $S = \int dt d^3x \mathcal{L}(\phi, \partial\phi/\partial x_j)$
- \mathcal{L} is built up from the field and its gradients
- Given \mathcal{L} we can recover the field eqns from the Euler-Lagrange eqns

Example: wave eqn

$$\frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} - \frac{\partial^2 \phi}{\partial x^2} = 0 \quad \leftrightarrow \quad \mathcal{L} = \underbrace{\frac{1}{c^2} \left(\frac{\partial \phi}{\partial t} \right)^2}_{\text{KE}} - \underbrace{\left(\frac{\partial \phi}{\partial x} \right)^2}_{\text{PE}}$$

- Wave eq linear in $\phi \quad \leftrightarrow \quad \mathcal{L}$ quadratic in ϕ
- Both d.e. and \mathcal{L} unchanged by Lorentz transformation

$$\begin{pmatrix} ct' \\ x' \end{pmatrix} = \begin{pmatrix} \gamma & \beta\gamma \\ \beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} ct \\ x \end{pmatrix} \quad \text{and} \quad \phi(x', t') = \phi(x, t)$$

Example: emag

- Now field has 4 components $A_x, A_y, A_z, A_t = -\phi/c$

$$B = \nabla \times A \quad \frac{E}{c} = \nabla A_0 - \frac{1}{c} \frac{\partial A}{\partial t}$$

↑ Electrostatics ↑ Faraday

$$\frac{1}{c^2} \frac{\partial^2 A_j}{\partial t^2} - \sum_{i=1}^3 \frac{\partial^2 A_j}{\partial x_i^2} = 0 \quad \leftrightarrow \quad \mathcal{L} = \frac{1}{2\mu_0 c} \left(\frac{E^2}{c^2} - B^2 \right)$$

↑ KE ↑ PE

- \mathcal{L} again quadratic in field gradients
- \mathcal{L} again unchanged by Lorentz transformation given that A_i changes by same rule as x_i (is a vector)

$$\begin{pmatrix} A'_t \\ A'_x \end{pmatrix} = \begin{pmatrix} \gamma & \beta\gamma \\ \beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} A_t \\ A_x \end{pmatrix}$$

Dirac

1902-1984



- Dirac was seeking a relativistic generalisation of the Schroedinger eq
- He discovered a new aspect of the aether – a four-component, complex-valued field $\psi = (\psi_1, \psi_2, \psi_3, \psi_4)$
- It transforms in a new way under Lorentz transformations – it's a 'Dirac spinor' $\psi \rightarrow \psi' = M(\beta)\psi$
- Electrons & positrons are excitations of ψ just as photons are excitations of A

Dirac equation

$$i \sum_j \gamma_j \frac{\partial \psi}{\partial x_j} - \frac{2\pi}{\lambda_C} \psi = 0 \quad \Leftrightarrow \quad \mathcal{L} = i \sum_j \bar{\psi} \gamma_j \frac{\partial \psi}{\partial x_j} - \frac{2\pi}{\lambda_C} \bar{\psi} \psi$$

\uparrow 4×4 matrix Compton λ

Dirac eqn

$$\gamma_t = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \quad \gamma_x = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$

- Linear coupled wave eqns now 1st order (hence i)
- \mathcal{L} unchanged by Lorentz transformation given that ψ changes as a spinor

Coupling fields

$$\frac{1}{c^2} \frac{\partial^2 A_j}{\partial t^2} - \sum_{i=1}^3 \frac{\partial^2 A_j}{\partial x_i^2} = -\mu_0 \underset{\substack{\uparrow \\ \text{antenna}}}{J_j} \quad \leftrightarrow \quad \mathcal{L} = \frac{1}{2\mu_0 c} \left(\frac{E^2}{c^2} - B^2 \right) + \frac{1}{c} \underset{\substack{\uparrow \\ \text{Current density}}}{J} \cdot A$$

- But current density from electrons is

$$J_i = e \bar{\psi} \gamma_i \psi$$

- So \mathcal{L} is

$$\mathcal{L} = \underbrace{\frac{1}{2\mu_0 c} \left(\frac{E^2}{c^2} - B^2 \right)}_{\text{Emag waves}} + \underbrace{\frac{e}{c} A_i \bar{\psi} \gamma_i \psi}_{\text{coupling}} + \underbrace{i \bar{\psi} \gamma_i \frac{\partial \psi}{\partial x_i} - \frac{2\pi}{\lambda_C} \bar{\psi} \psi}_{\text{Free electrons}}$$

- Coupling term is here cubic in the fields

Summary

- The spacetime (aether) is a complex dynamical system that carries several (many-component) fields
- We can derive all the wave eqns that describe this system from a single Lagrangian density
- Lorentz invariance of L **strongly** restricts the possible field eqns!
- It comprises a block for each mode (free particles) plus coupling terms between blocks