

Black holes in the nearby universe

John Magorrian

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Different “types” of black hole in astrophysics

Approx. 10 002 black holes around the Andromeda Galaxy.



Overview

(At least) two different “types” of astrophysical BHs.

	Stellar mass	Supermassive	
Mass [M_{\odot}]	10	10^6 – 10^9	$M_{\odot} = 2 \times 10^{30}$ kg
Horizon	30 km	\sim AU	
Where	among stars	galaxy centres	
Origin	dead star	???	
θ_{horizon} ["]	10^{-5}	10^{-5}	$3600'' = 1^{\circ}$

Optical/IR telescopes resolve $\Delta\theta \sim 0''.1 \gg \theta_{\text{horizon}}$:

- BH looks like Newtonian point mass, $\Phi = -GM/r$;
- can hope to measure M , but not spin (at least not directly);

Find compact M , then argue that it *must* be a BH.

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- 1 Stellar-mass black holes
 - one example
- 2 “Supermassive” black holes
 - three examples
 - consequences

I. Stellar-mass BHs

(Remnants of dead stars)

Q: What stops stars from collapsing? **A:** Pressure.

- thermal pressure in normal stars;
- electron degeneracy pressure in white dwarfs;
- neutron degeneracy pressure in neutron stars.

There is a maximum mass for the latter.

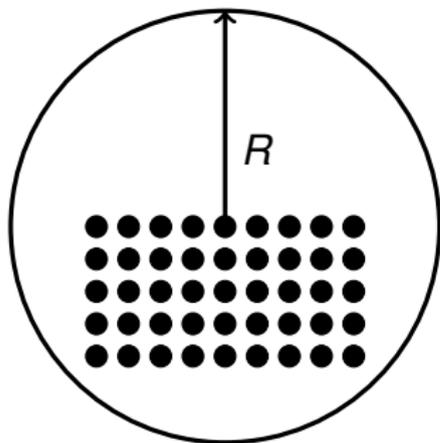
Identical fermions refuse to share states with one another. When $T \rightarrow 0$, states up to some maximum energy E_F , the **Fermi energy**, are fully occupied.

For n ultrarelativistic fermions/unit volume, $E_F \sim n^{1/3} \hbar c$.

Neutron stars

Landau's (1932) derivation of maximum mass

Assume: equilibrium = state of minimum total energy.



N neutrons, mass m ,
within radius R .

Fermi energy is

$$\varepsilon_F \sim \hbar c \left(\frac{N}{R^3} \right)^{1/3}.$$

(ultrarel. degenerate fermions)

Star's total energy

$$\begin{aligned} E(N, R) &\sim N\varepsilon_F - \frac{GM^2}{R} \\ &= \frac{\hbar c N^{4/3} - G(mN)^2}{R}. \end{aligned}$$

Numerator $< 0 \Rightarrow E \downarrow$ if $R \downarrow$: collapse!

So, for stability we must have

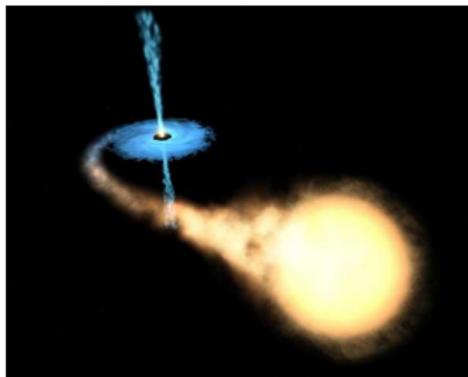
$$M = Nm \lesssim \left(\frac{\hbar c}{Gm^2} \right)^{3/2} m.$$

More complete models with GR and $P(\rho)$ give $M \simeq 3.1 M_\odot$.

Observations of X-ray binaries

Circular binary with masses M_{\bullet} and M_{\star} , separation r .

We observe \star , not \bullet .



Binary's orbital period

$$T = 2\pi \sqrt{\frac{r^3}{G(M_{\bullet} + M_{\star})}}$$

Peak line-of-sight velocity

$$v_{\star} = \frac{M_{\bullet}}{M_{\bullet} + M_{\star}} v \sin i,$$

where

$$v = \sqrt{\frac{G(M_{\bullet} + M_{\star})}{r}}$$

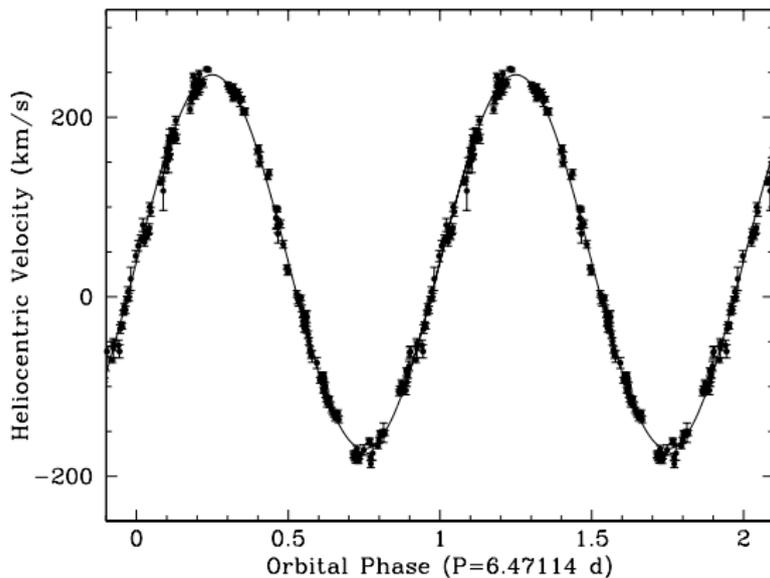
Therefore the quantity

$$f(M_{\bullet}) \equiv \frac{T v_{\star}^3}{2\pi G} = \frac{M_{\bullet} \sin^3 i}{(1 + M_{\star}/M_{\bullet})^2} < M_{\bullet}.$$

Example: V404 Cygni

(Shahbaz et al. 1994)

V404 Cygni: variable star/nova/X-ray transient, ~ 2 kpc away.



$$v_* \sin i \simeq 220 \text{ km/s}$$

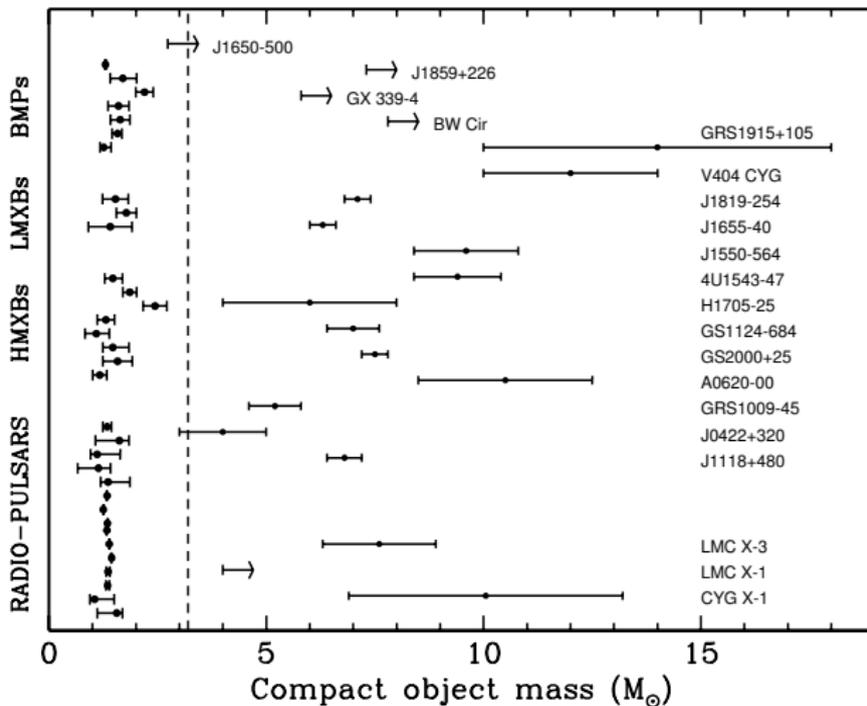
$$T = 6.47 \text{ d}$$

$$\Rightarrow f(M) \simeq 7 M_\odot$$

Detailed modelling
constrains also i
and M_* .

Demographics of stellar remnants

(Casares 2006)



Upper bound $f(M_{\bullet})$ shrunk by detailed modelling of binary system to constrain (i, M_{\star}) .

II. Supermassive black holes

(origin unknown)

Why expect supermassive BHs?

Observational: QSOs (Quasi-Stellar objects), 1960's

- Distant ($z \gtrsim 0.1$) point-like sources;
- power outputs $\sim 10^{12} L_{\odot} \sim 10^{39} \text{ W}$;
- (some) variable on \sim day timescales
 - no more than light days across

Interpretation:

- Lynden-Bell (1969): QSOs caused by accretion by “Schwarzschild throats” located in the middle of galaxies;
- Rees (1978, 1984): supermassive BHs should be easy to form.

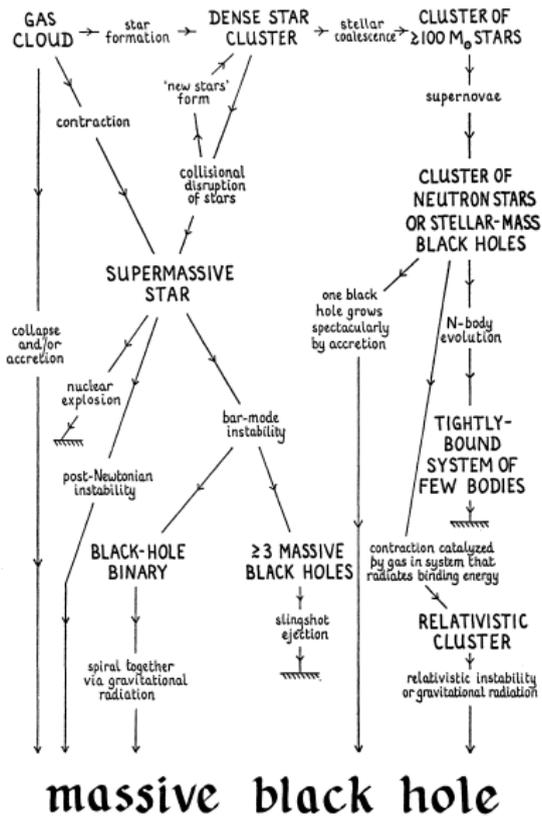
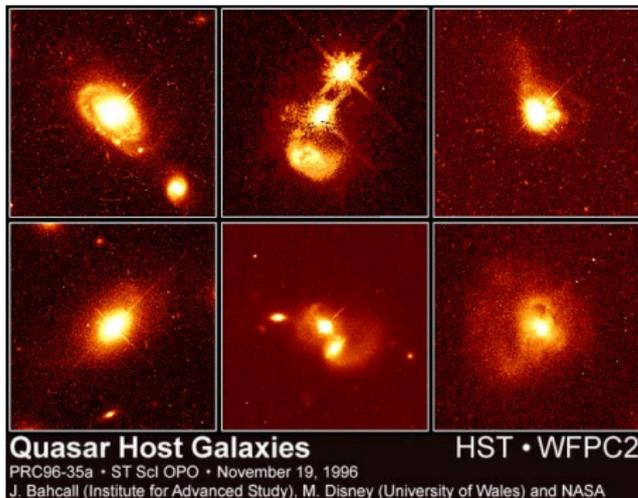


FIG. 2

Possible modes of formation of a massive black hole in a galactic nucleus.

Nowadays: less certain.

Why expect supermassive BHs?



1996

Hubble Space Telescope confirms that QSOs are located in galaxies.

Co-moving number density has decreased ~ 100 -fold between $z = 2$ and present.

Where are the dead QSOs now?

Use dynamical tracers (stars, gas) to probe central potentials of nearby, inactive galaxies.

Orbits of stars around BH

(A reminder)

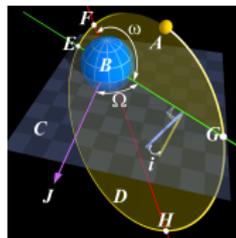
Stars move in gravitational potential

$$\Phi(\mathbf{r}) = -\frac{GM_{\bullet}}{r} + \Phi_{\star}(\mathbf{r}).$$

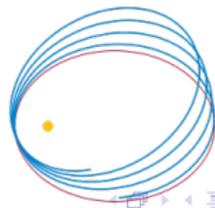
Close to BH, Φ_{\star} is small perturbation: akin to solar system!

To 0th-order ($\Phi_{\star} \rightarrow 0$), orbits are closed ellipses with constant

- semimajor axis a ,
- eccentricity e ,
- argument of periapse ω ,
- angles (i, Ω) giving orientation of orbit plane / \mathbf{L} vector.



To next order ($\Phi_{\star} \neq 0$), ω **precesses** with frequency $\propto M_{\star}$



The Milky Way

(Genzel et al.; Ghez et al.)

We're ~ 8 kpc from Galactic centre.

Undetectable in optical: dust attenuates by factor $\sim 10^{-10}$.

Radio:

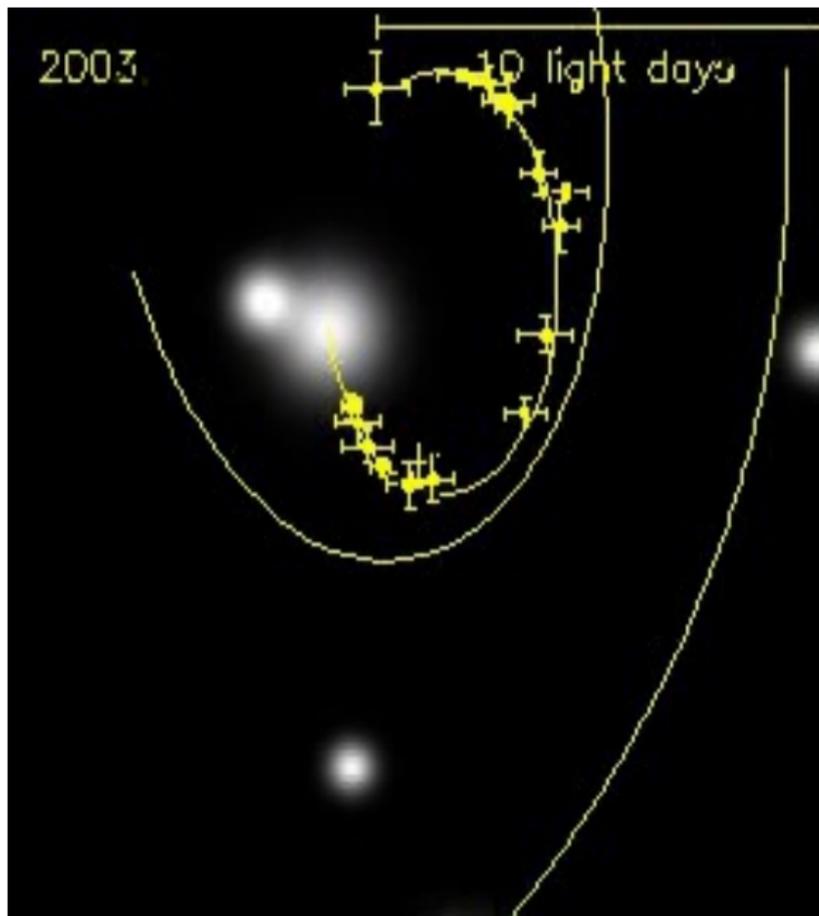
Compact source Sgr A* (1974).

Size 37 microarcsec (VLBI observations, 2008): ~ 1 AU.

Infrared: (1992–present)

2 groups start long-term monitoring.

See stars at GC move!

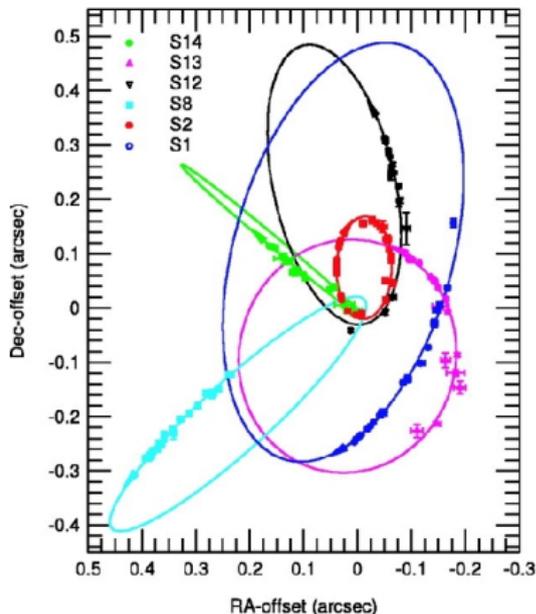
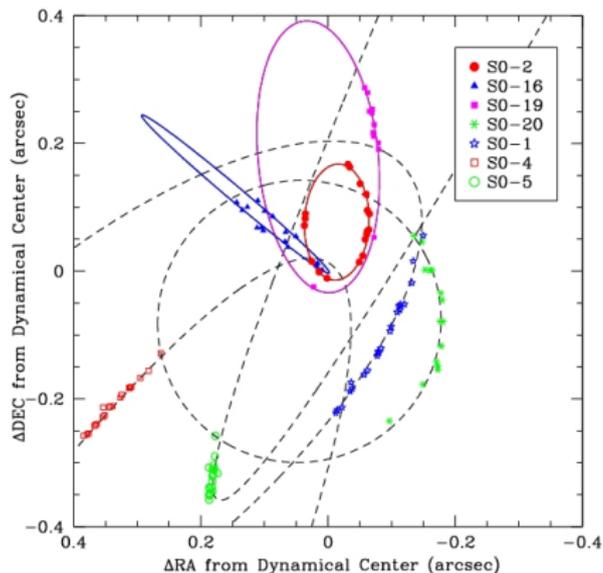


Movie credit: MPE

The massive dark object at the Galactic Centre

(Genzel+MPE, Ghez+UCLA)

Orbits of *S stars* around Sgr A*:



All fit by Kepler ellipses around $M_{\bullet} = 4.4 \times 10^6 M_{\odot}$ at Sgr A*.

NB: star S2 has period 15.2 yr, pericenter 17 light hours.

Is this massive dark object a BH?

Lack of detectable deviations from Keplerian ellipse for S2 \Rightarrow
extent of dark object $\lesssim 10^{-3}$ pc.

Suppose object is cluster of things of mass m_* spread over
radius $r_{1/2}$.

Cluster “evaporates” on timescale (Maoz 1998)

$$t_{\text{evap}} = 10^8 \left(\frac{M_\bullet}{4 \times 10^6 M_\odot} \right)^{1/2} \left(\frac{M_\odot}{m_*} \right) \left(\frac{r_{1/2}}{0.01 \text{ pc}} \right)^{3/2} \text{ yr.}$$

This rules out any plausible, known cluster constituents.

Maybe it's something more exotic than BH?

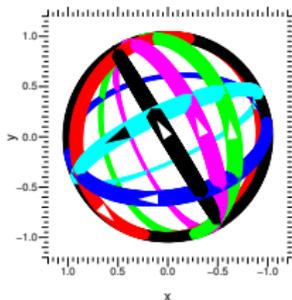
Puzzles among the tracers

Naive expectation:

- only old stars
- completely relaxed, Maxwell-Boltzmann $f(\mathbf{p})$

Reality:

- 1 Dominant old, relaxed population
 - Observations sample only the very brightest stars
- 2 S stars:
 - B stars with masses $3.5 - 20 M_{\odot}$
 - found at radii $r \lesssim 0.02$ pc
 - randomized orbits
- 3 Population of ~ 100 massive young stars in innermost pc.
 - O/WR stars, born with $30 - 100 M_{\odot}$.
 - Arranged in two eccentric, warped discs.
 - Discs are misaligned with each other,
 - but have identical 6 ± 2 Myr ages.



(Paumard+2006: Black, blue, red)

Warped discs around a BH

(Rauch & Tremaine 1997; Tremaine 2005)

When Φ_* is small, $t_{\text{orb}} \ll t_{\text{precess}} \sim t_{\text{orb}}(M_\bullet/M_*)$.

On timescales $\lesssim t_{\text{precess}}$, replace stars by orbit-averaged elliptical rings.

Each ring labelled by $(a, e, i, \omega, \Omega)$.

Elliptical rings exert torques on each other (Laplace–Lagrange).
Orientation of each ring undergoes a random walk.

Warps, misalignments develop on timescales Nt_{precess} , where $N = M_*/m_*$ is the number of rings.

At $r \simeq 0.5$ pc from GC BH, $t_{\text{precess}} \sim 10^6 M_\odot$:
consistent with 6 Myr stellar ages.

How long before the BH consumes us?

Stars of mass m_* , radius r_* that come within radius

$$r_t \approx \left(\frac{M_\bullet}{m_*} \right)^{1/3} r_*$$

of the BH are *tidally disrupted*.

$\sim \frac{1}{2}$ debris swallowed by BH.

$r_t \sim 10r_S$ for sun at GC.



Animation: ESA

Only stars with angular momenta $L^2 \lesssim 2GMr_t$ have $r_{\text{peri}} < r_t$.

Total mass of such stars in MW: $\sim 10 M_\odot$!

Composed of plunging radial orbits, unbound to BH.

Fluctuations in Φ_* give stars small kicks in L^2 .

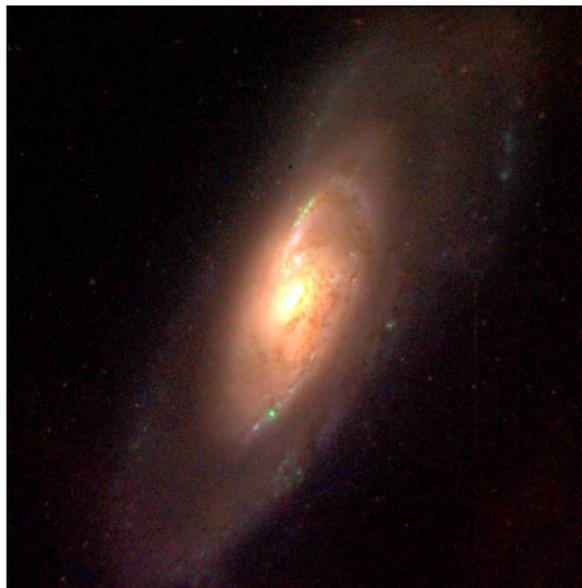
Steady-state delivery rate $\dot{M}_\bullet \sim 10^{-4} M_\odot/\text{yr}$.

\Rightarrow **BHs don't grow by eating stars!**

NGC 4258

(aka Messier 106, ~ 10 Mpc distant)

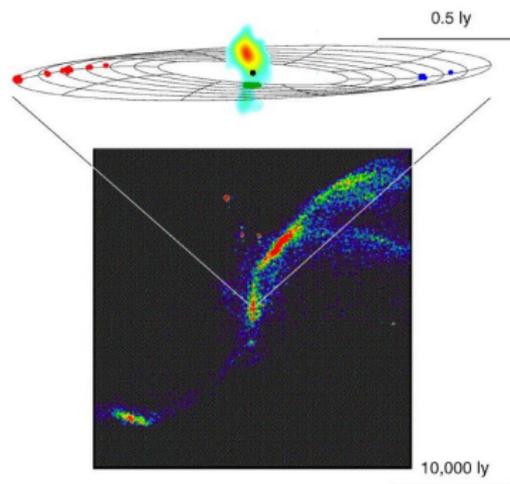
Optical image:



NGC 4258

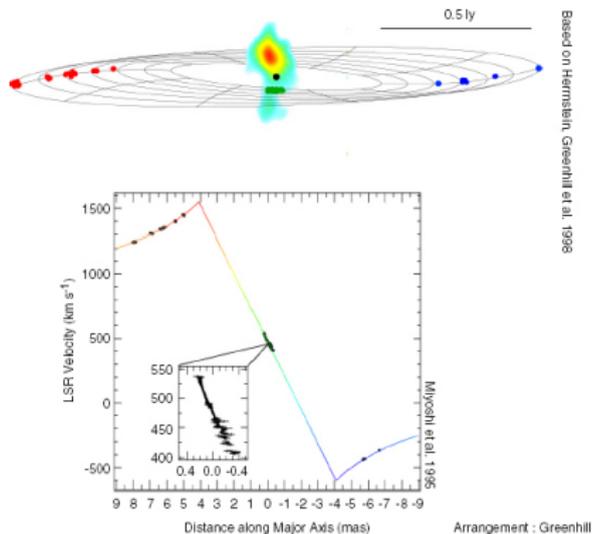
(Miyoshi et al. 1995)

Radio interferometry reveals a central disc of molecular gas, with water maser emission.



Beautifully Keplerian. $M_{\bullet} = 4 \times 10^7 M_{\odot}$.

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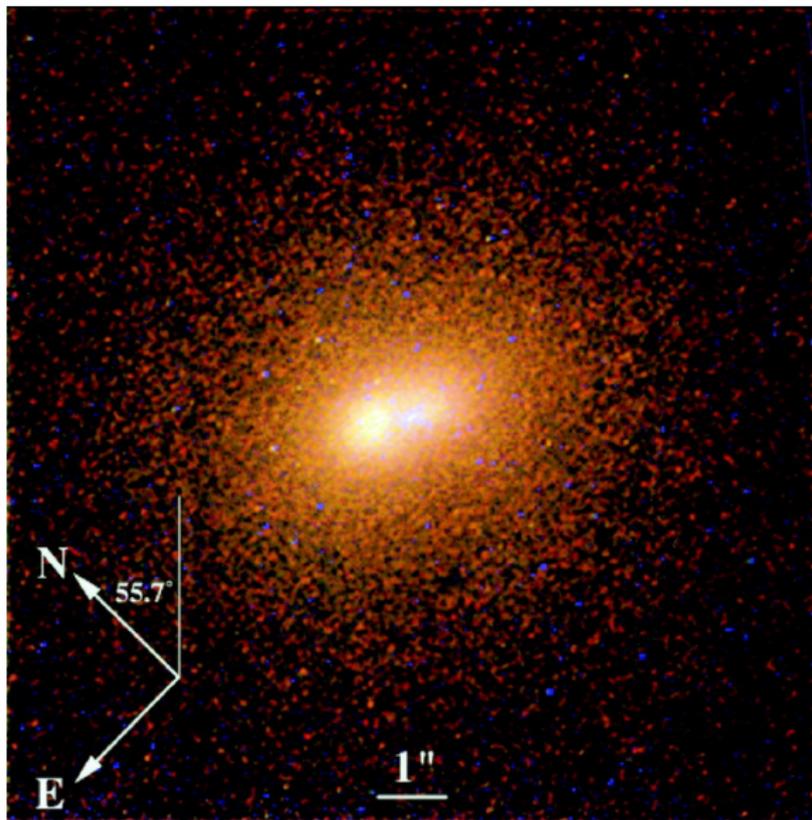


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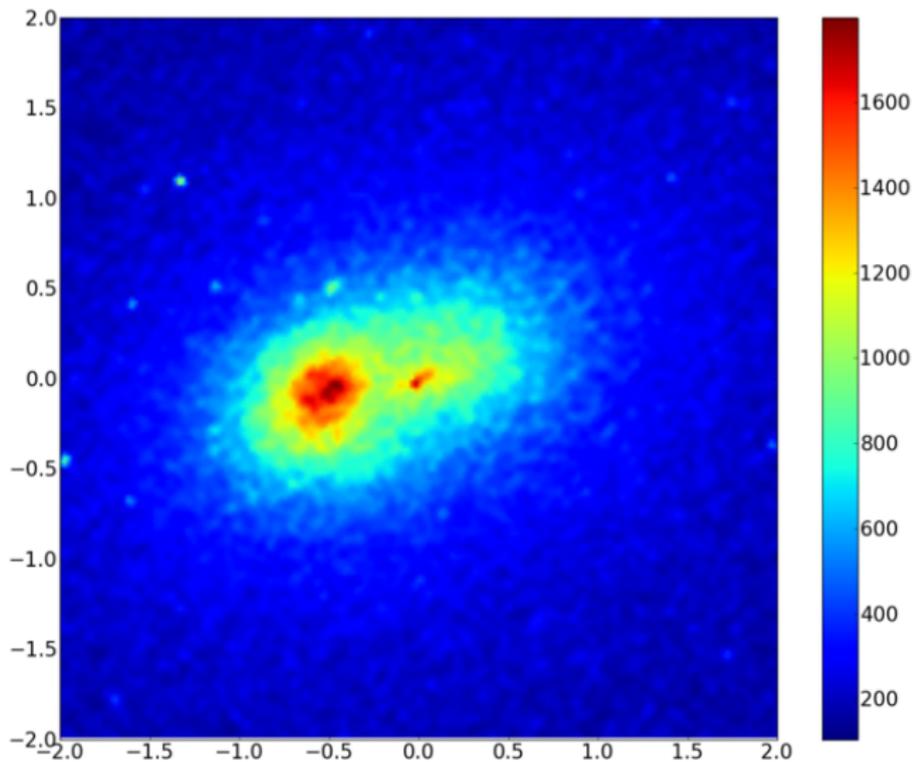
The black hole at the centre of Andromeda (M31)



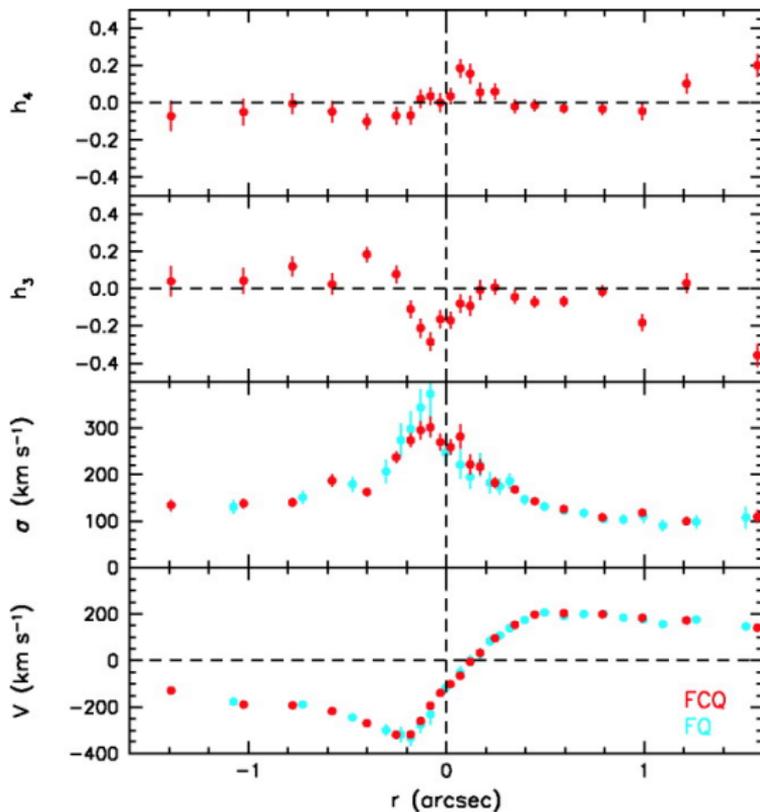
True-colour HST photometry (Lauer et al 98):



Single-band false colour image (Lauer et al 98):



STIS CaT long-slit kinematics (Bender et al 2005):



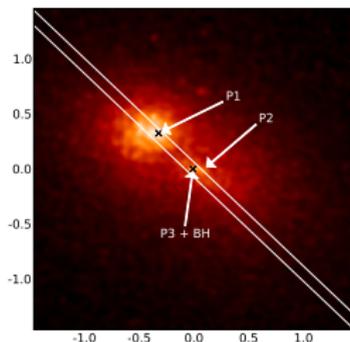
Characteristic numbers for M31 double nucleus

Distinct nucleus $L_V \simeq 6 \times 10^6 L_\odot$ with two peaks, P1 and P2.

P1 and P2 have identical colours: **old, red**.

P2 is photometric centre of galaxy (to $\sim 0.1''$).

- P1-P2 separation $r = 0.5$ arcsec = 2 parsec
- $\Delta v \sim 200$ km/s
- Dynamical time $2\pi r/v \sim 10^5$ yr.



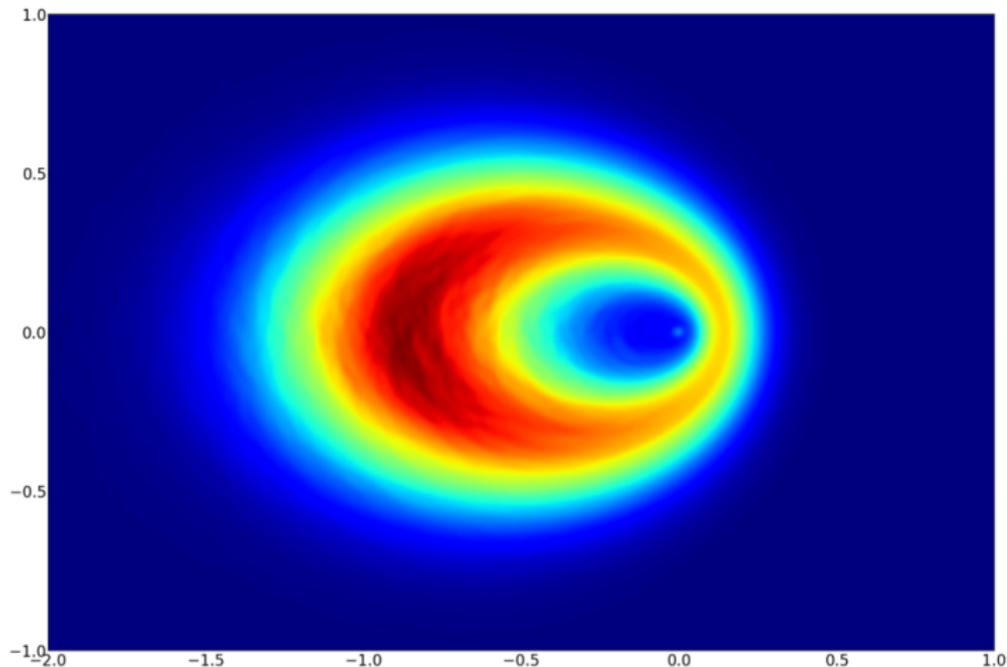
Two distinct clusters?

Dynamical friction timescale
only $\sim 10^8$ yr.

An eccentric disc around a BH?

Tremaine 1995

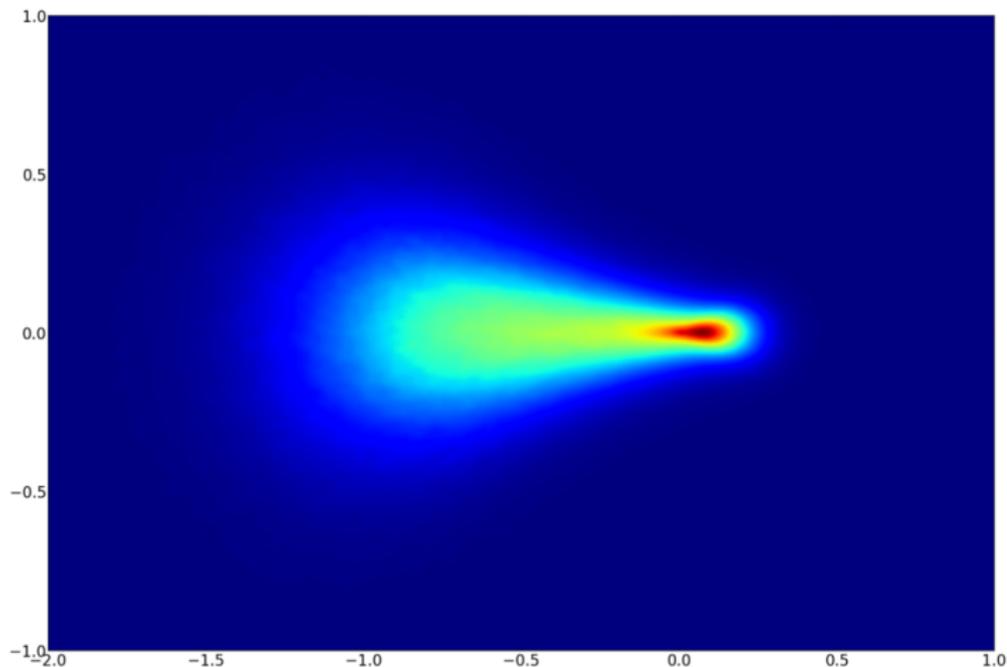
Clump of orbits around $a = 1$, $e = 0.7$ viewed top down:



An eccentric disc around a BH?

Tremaine 1995

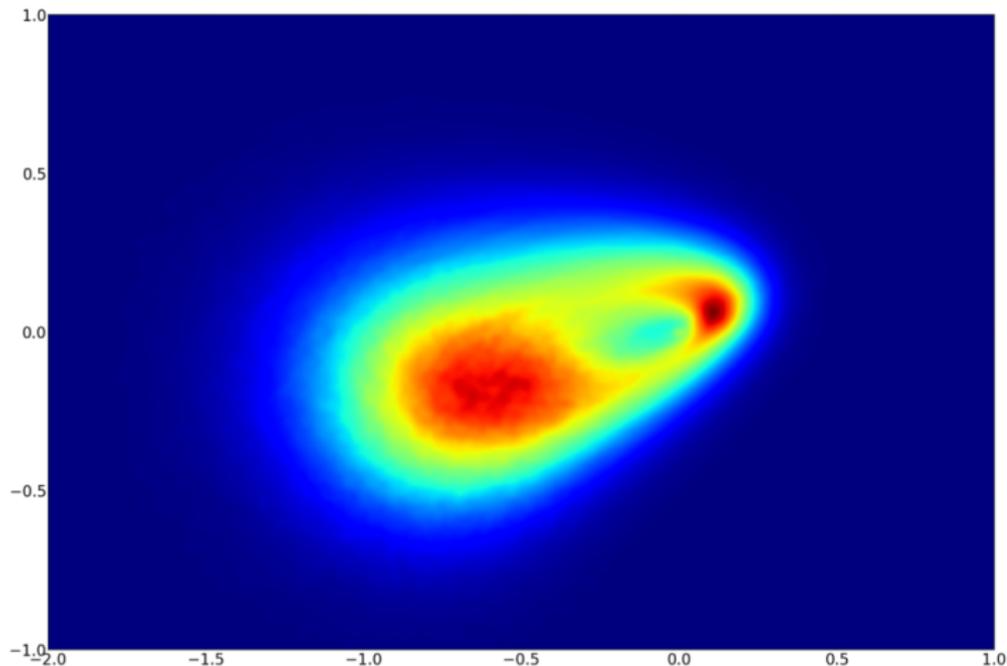
Clump of orbits around $a = 1$, $e = 0.7$ viewed edge on:



An eccentric disc around a BH?

Tremaine 1995

Clump of orbits around $a = 1$, $e = 0.7$ viewed just right:

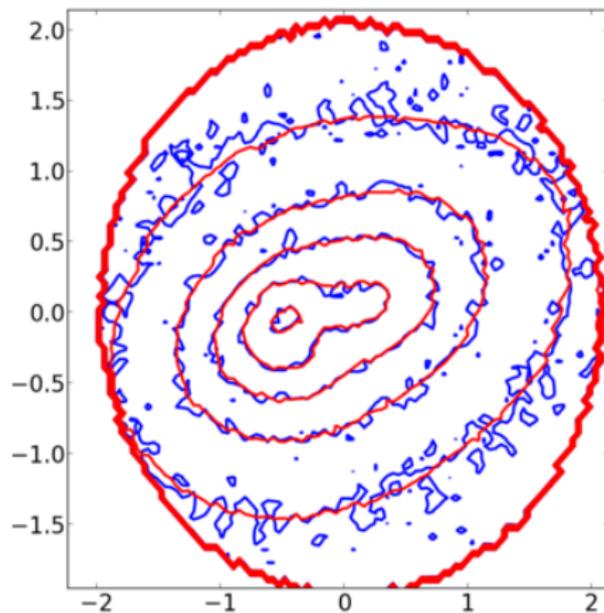


Naive models of M31 nucleus

(Brown & JM, 2013, 2015)

Assume biaxial symmetry in y and z planes (no warp!).
Find distribution $f(a, e, i, \omega, \Omega)$ that reproduces observations.

Results: Fit to photometry

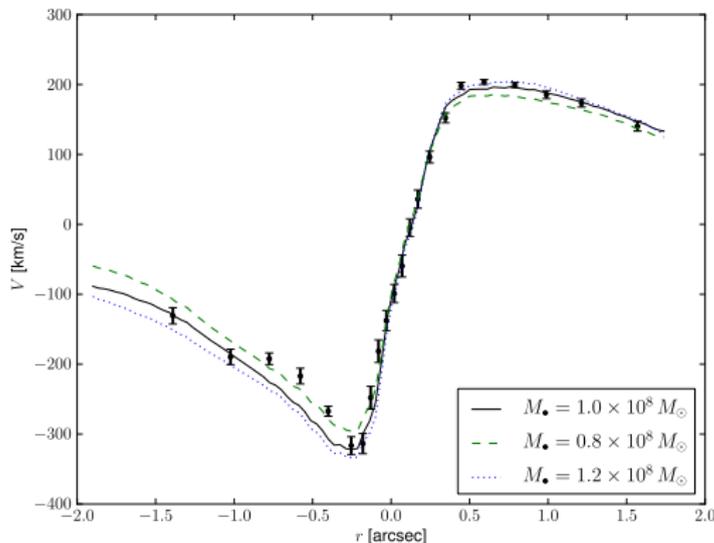


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Results: Fit to HST V for various assumed M_{\bullet} .

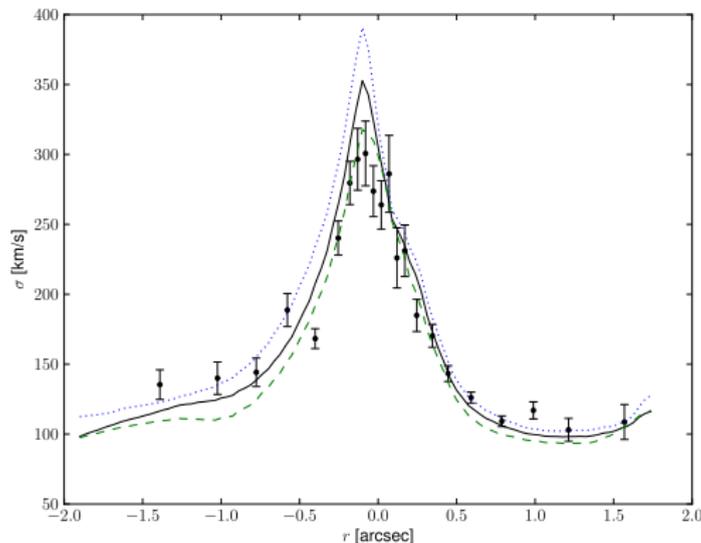


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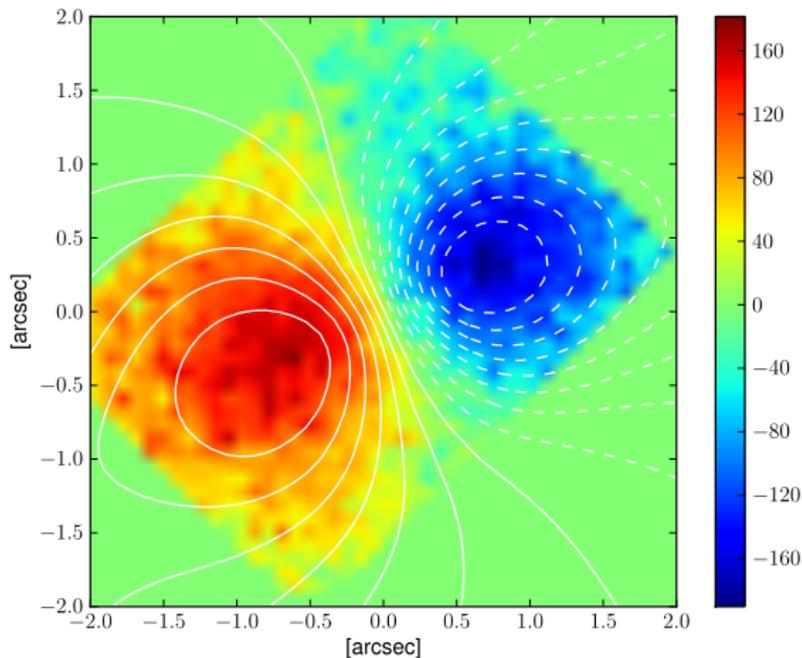
Assume biaxial symmetry in y and z planes (no warp!).
Find distribution $f(a, e, i, \omega, \Omega)$ that reproduces observations.

Results: Fit to HST σ for various assumed M_\bullet



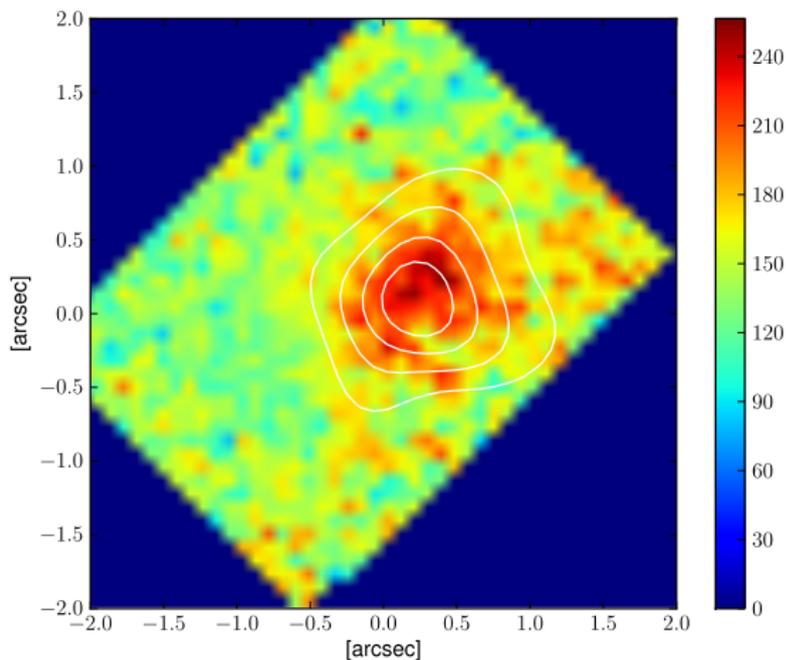
Naive models of M31 nucleus

Model *predictions* for $M_{\bullet} \simeq 10^8 M_{\odot}$ vs 2d kinematic maps: V



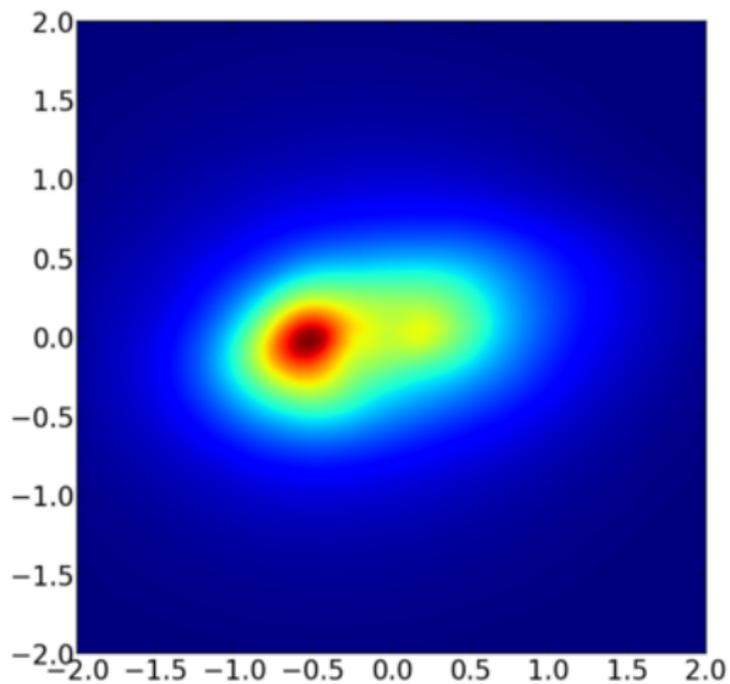
Naive models of M31 nucleus

Model *predictions* for $M_{\bullet} \simeq 10^8 M_{\odot}$ vs 2d kinematic maps: σ



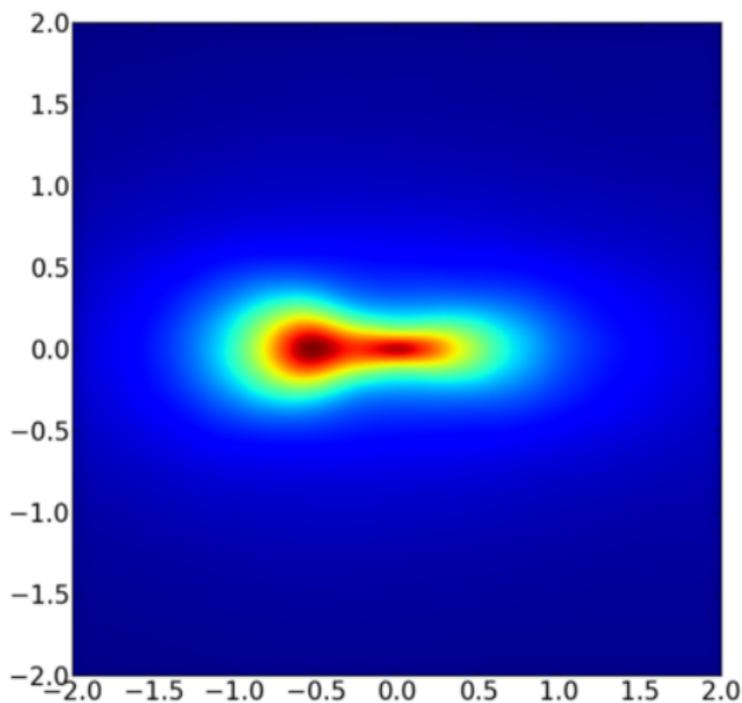
Naive models models of M31 nucleus

What does the disc look like? LOS projection:



Naive models of M31 nucleus

What does the disc look like? Edge-on:



Naive models: a puzzle

When we examine our best-fit $f(a, e, i, \omega, \Omega)$, we find

$$\frac{\langle i^2 \rangle}{\langle e^2 \rangle} \simeq 2.$$

But standard dynamical models of two-body relaxation in *circular* discs predict

$$\frac{\langle i^2 \rangle}{\langle e^2 \rangle} = \frac{1}{2}.$$

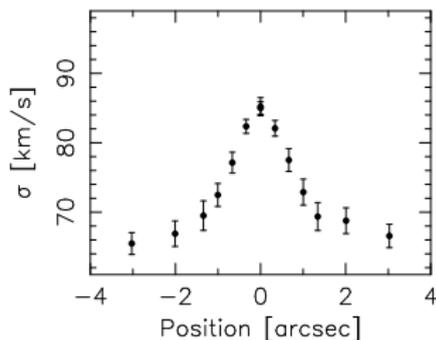
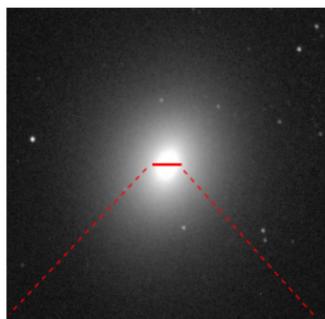
Our result seems robust, subject to biaxial assumption.

III. Supermassive BHs are standard equipment in galaxy bulges

Other galaxies

Q:What stops galaxy from collapsing? **A:** pressure!

Example: kinematics at very centre of M32

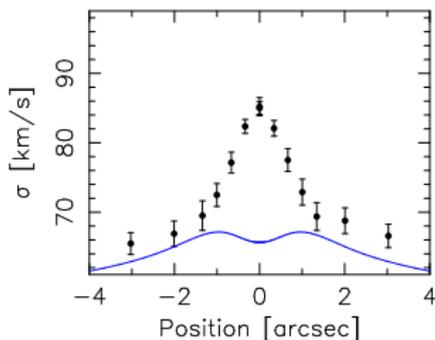
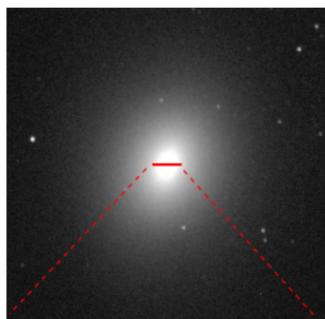


Must resolve radii where BH dominates Φ : nearby galaxies only.

Other galaxies

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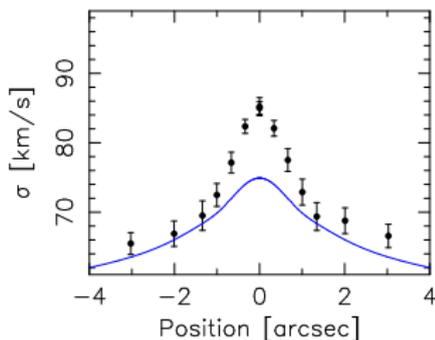
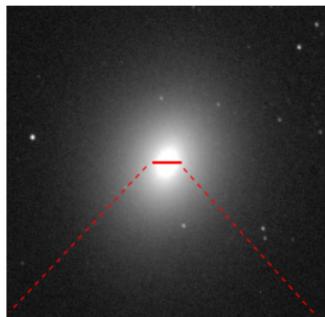
Model with no BH

Must resolve radii where BH dominates Φ : nearby galaxies only.

Other galaxies

Q:What stops galaxy from collapsing? **A:** pressure!

Example: kinematics at very centre of M32



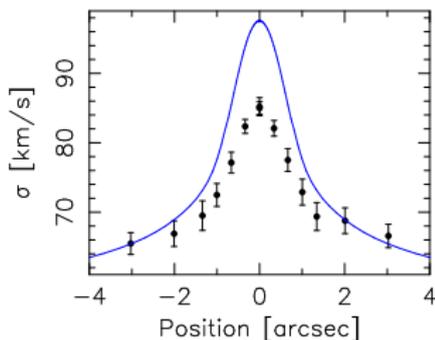
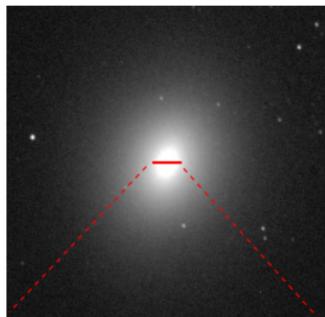
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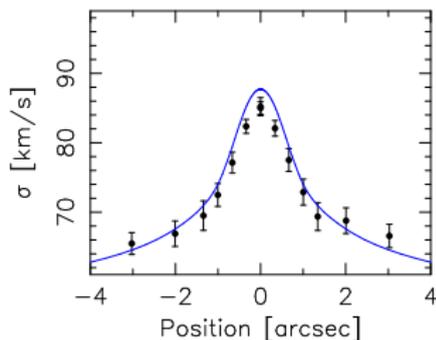
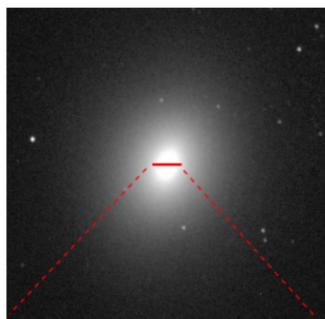
$$M_{\bullet} = 4 \times 10^6 M_{\odot}$$

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Other galaxies

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Example: kinematics at very centre of M32



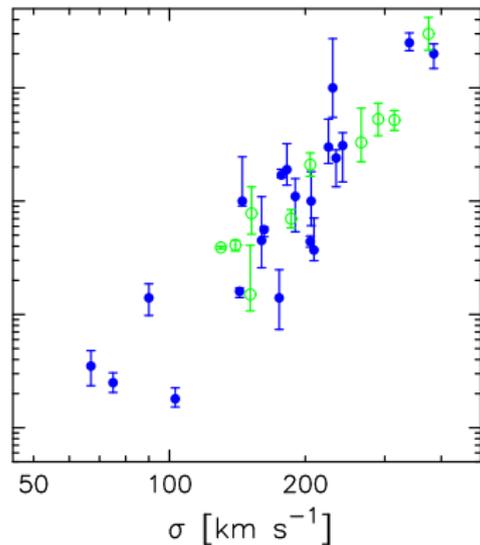
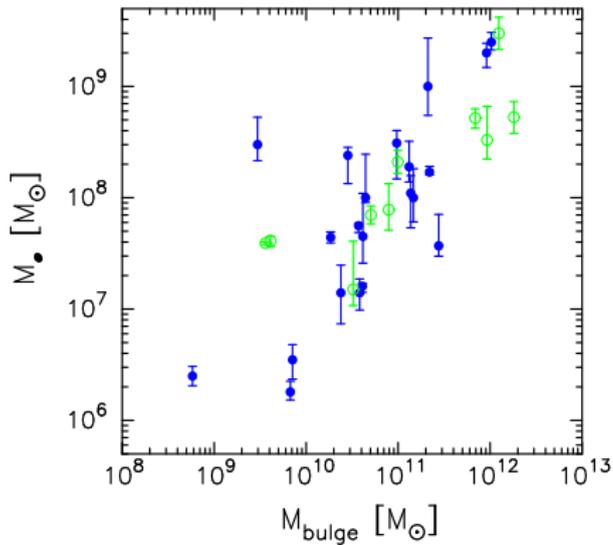
$$M_{\bullet} = 2.6 \times 10^6 M_{\odot}$$

Must resolve radii where BH dominates Φ : nearby galaxies only.

The most remarkable fact about BHs

(Gebhardt et al; Ferrarese & Merritt; Tremaine et al.)

BHs know about their host bulges:



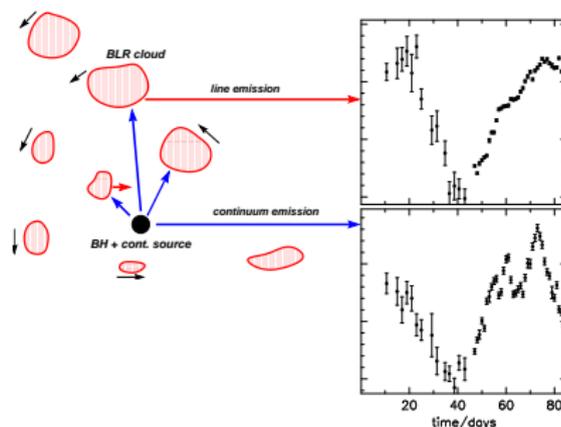
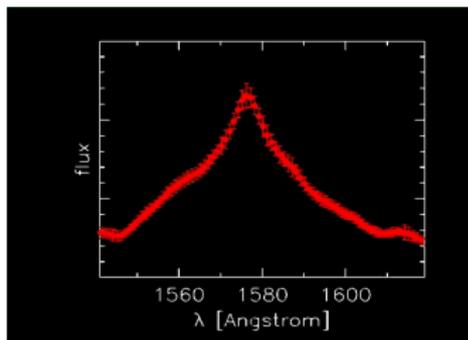
Upper left definitely empty (ish); bottom right almost certainly so.

$$M_{\bullet} \approx 10^{-3} M_{\text{bulge}}; \quad M_{\bullet} \sim \sigma^4.$$

Reverberation mapping

We can measure M_{\bullet} directly only in nearby, inactive galaxies.
What about distant, active ones? Broad Line Regions.

Time variation of CIV line in NGC 5548:



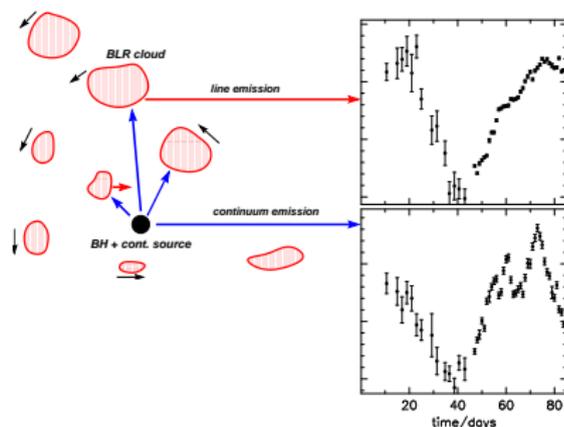
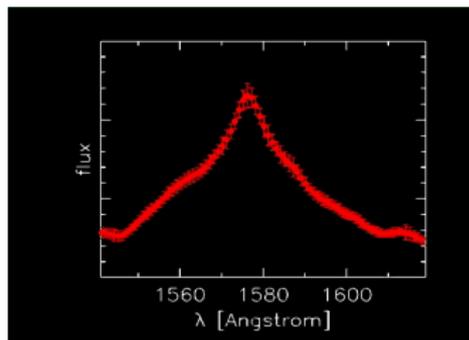
$M_{\bullet} = f r v^2 / G$, where v is line width, r is lag between line and continuum, and fudge factor f depends on geometry of BLR.

Works at very high redshift.

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Works at very high redshift.

- 1 Astrophysical evidence for BHs is compelling, but indirect:
 - limited by spatial resolution
 - No clear signature of event horizon (therefore spin)
 - but we can still measure masses.
 - At least two different ways of forming BHs.
 - In the future, gravitational-wave detectors should detect mergers of stellar-mass BHs.
- 2 Interesting dynamics in almost Keplerian potentials.
- 3 Most nearby galaxy bulges harbour a supermassive black hole
 - we don't understand how either form
 - but galaxy growth somehow linked to BH growth.